

**Numerical Methods of
Electromagnetic Field Theory I (NFT I)**
**Numerische Methoden der
Elektromagnetischen Feldtheorie I (NFT I) /**

5th Lecture / 5. Vorlesung

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3-D Electromagnetic Wave Propagation /
3D elektromagnetische Wellenausbreitung

Constitutive Equations for Vacuum /
 Konstituierende Gleichungen
 (Materialgleichungen) für Vakuum

$$\underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

Maxwell's equations / Maxwellsche Gleichungen

$$\left\{ \begin{array}{l} \frac{\partial}{\partial t} \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) \\ \frac{\partial}{\partial t} \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) \\ \nabla \cdot \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \rho_m(\underline{\mathbf{R}}, t) \\ \nabla \cdot \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \rho_e(\underline{\mathbf{R}}, t) \end{array} \right.$$

Continuity equations / Kontinuitätsgleichungen

$$\nabla \cdot \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_e(\underline{\mathbf{R}}, t)$$

The diagram shows a flow from the left side where constitutive equations for vacuum are listed. An arrow points down to a central point, which then branches into two arrows pointing to the right side. The right side contains Maxwell's equations and continuity equations, each preceded by a brace indicating they are grouped together.

$$\frac{\partial}{\partial t} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0} \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

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$$\frac{\partial}{\partial t} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (1)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0} \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (2)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \frac{\partial}{\partial t} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (3)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0} \nabla \times \frac{\partial}{\partial t} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (4)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\mu_0} \nabla \times \left[-\frac{1}{\epsilon_0} \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \right] - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (5)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0} \nabla \times \left[-\frac{1}{\mu_0} \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \right] - \frac{1}{\epsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (6)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) + \frac{1}{\epsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (7)$$

$$\frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (8)$$

3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\frac{1}{\epsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) + \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = +\frac{1}{\epsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) - \frac{1}{\mu_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (1)$$

$$\frac{1}{\epsilon_0 \mu_0} \nabla \times \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) + \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = -\frac{1}{\epsilon_0 \mu_0} \nabla \times \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) - \frac{1}{\epsilon_0} \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (2)$$

$$c_0 = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

$$-\nabla \times \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (3)$$

$$-\nabla \times \nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (4)$$

Vector identity /
Vektoridentität $\nabla \times \nabla \times = \nabla \nabla \bullet - \underbrace{\nabla \bullet \nabla}_{=\nabla^2 = \Delta} = \nabla \nabla \bullet - \Delta$ Short-hand notation /
Abkürzende
Schreibweise $\nabla \bullet \nabla = \nabla^2 = \Delta$

$$-[\nabla \nabla \bullet - \Delta] \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) \quad (5)$$

$$-[\nabla \nabla \bullet - \Delta] \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_{\text{m}}(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_{\text{e}}(\underline{\mathbf{R}}, t) \quad (6)$$

3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$-\left[\nabla \nabla \bullet - \Delta\right] \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$-\left[\nabla \nabla \bullet - \Delta\right] \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

3rd and 4th Maxwell's
equations / 3. und 4.
Maxwellsche Gleichung

$$\nabla \bullet \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \bullet \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \rho_e(\underline{\mathbf{R}}, t)$$

Constitutive equations /
Materialgleichungen

$$\underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

$$\nabla \bullet \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = \frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \bullet \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \frac{1}{\epsilon_0} \rho_e(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \nabla \underbrace{\nabla \bullet \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)}_{=\frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}}, t)} - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \nabla \underbrace{\nabla \bullet \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)}_{=\frac{1}{\epsilon_0} \rho_e(\underline{\mathbf{R}}, t)} - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \nabla \left[\frac{1}{\mu_0} \rho_m(\underline{\mathbf{R}}, t) \right] - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \nabla \left[\frac{1}{\epsilon_0} \rho_e(\underline{\mathbf{R}}, t) \right] - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \frac{1}{\mu_0} \nabla \rho_m(\underline{\mathbf{R}}, t)$$

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \frac{1}{\epsilon_0} \nabla \rho_e(\underline{\mathbf{R}}, t)$$

Laplace operator in Cartesian coordinates /
Laplace-Operator in Kartesischen Koordinaten

$$\Delta = \nabla \bullet \nabla$$

$$= \left(\mathbf{e}_x \frac{\partial}{\partial x} + \mathbf{e}_y \frac{\partial}{\partial y} + \mathbf{e}_z \frac{\partial}{\partial z} \right) \bullet \left(\mathbf{e}_x \frac{\partial}{\partial x} + \mathbf{e}_y \frac{\partial}{\partial y} + \mathbf{e}_z \frac{\partial}{\partial z} \right)$$

$$= \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \cdot [\nabla \mathbf{E}(\underline{\mathbf{R}}, t)] \\ = \left(\underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot \left[\left(\underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \mathbf{E}(\underline{\mathbf{R}}, t) \right]$$

$$\begin{array}{ll} \text{Short-hand notation /} & \frac{\partial}{\partial x} = \partial_x \\ \text{Abkürzende Schreibweise} & \frac{\partial}{\partial y} = \partial_y \quad \frac{\partial}{\partial z} = \partial_z \quad \frac{\partial}{\partial t} = \partial_t \end{array}$$

$$\begin{aligned}\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)] \\ &= (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \cdot \left[(\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) \right]\end{aligned}$$

$$\begin{aligned} & \left(\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z \right) \left[E_x(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_x + E_y(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_y + E_z(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_z \right] \\ &= \underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\underline{\mathbf{R}}, t) \\ &+ \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_y E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\underline{\mathbf{R}}, t) \\ &+ \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_z E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_z E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\underline{\mathbf{R}}, t) \end{aligned}$$

3-D Electromagnetic Wave Propagation /
3D elektromagnetische Wellenausbreitung

$$\begin{aligned}\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)] \\&= (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \cdot \left[(\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) \right] \\&= (\underline{\mathbf{e}}_x \partial_x + \underline{\mathbf{e}}_y \partial_y + \underline{\mathbf{e}}_z \partial_z) \cdot \left[\begin{aligned} &\underline{\mathbf{e}}_x \underline{\mathbf{e}}_x \partial_x E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_y \partial_x E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_x \underline{\mathbf{e}}_z \partial_x E_z(\underline{\mathbf{R}}, t) \\&+ \underline{\mathbf{e}}_y \underline{\mathbf{e}}_x \partial_y E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_y \partial_y E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \underline{\mathbf{e}}_z \partial_y E_z(\underline{\mathbf{R}}, t) \\&+ \underline{\mathbf{e}}_z \underline{\mathbf{e}}_x \partial_z E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_y \partial_z E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \underline{\mathbf{e}}_z \partial_z E_z(\underline{\mathbf{R}}, t) \end{aligned} \right]\end{aligned}$$

3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\begin{aligned}
\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \nabla \cdot [\nabla \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)] \\
&= \left[\underline{\mathbf{e}}_x \partial_x^2 E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \partial_y^2 E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \partial_z^2 E_z(\underline{\mathbf{R}}, t) \right] \\
&\quad + \left[\underline{\mathbf{e}}_x \partial_y^2 E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \partial_x^2 E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \partial_y^2 E_z(\underline{\mathbf{R}}, t) \right] \\
&\quad + \left[\underline{\mathbf{e}}_x \partial_z^2 E_x(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_y \partial_z^2 E_y(\underline{\mathbf{R}}, t) + \underline{\mathbf{e}}_z \partial_z^2 E_z(\underline{\mathbf{R}}, t) \right] \\
&= (\partial_x^2 + \partial_y^2 + \partial_z^2) \underbrace{[E_x(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_x + E_y(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_y + E_z(\underline{\mathbf{R}}, t) \underline{\mathbf{e}}_z]}_{=\underline{\mathbf{E}}(\underline{\mathbf{R}}, t)} \\
&= (\partial_x^2 + \partial_y^2 + \partial_z^2) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)
\end{aligned}$$

$$\begin{aligned}
\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= (\partial_x^2 + \partial_y^2 + \partial_z^2) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) \\
&= \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)
\end{aligned}$$

3-D Electromagnetic Wave Propagation / 3D elektromagnetische Wellenausbreitung

$$\begin{aligned}
\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) &= -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \frac{1}{\mu_0} \nabla \rho_m(\underline{\mathbf{R}}, t) \\
\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \frac{1}{\epsilon_0} \nabla \rho_e(\underline{\mathbf{R}}, t)
\end{aligned}$$

$$\begin{aligned}
\Delta \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= (\partial_x^2 + \partial_y^2 + \partial_z^2) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) \\
\Delta \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) &= (\partial_x^2 + \partial_y^2 + \partial_z^2) \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)
\end{aligned}$$

$$\begin{aligned}
\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) &= -\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \frac{1}{\mu_0} \nabla \rho_m(\underline{\mathbf{R}}, t) \\
\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) &= \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) + \frac{1}{\epsilon_0} \nabla \rho_e(\underline{\mathbf{R}}, t) \\
\nabla \cdot \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) &= -\frac{\partial}{\partial t} \rho_m(\underline{\mathbf{R}}, t) \\
\nabla \cdot \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) &= -\frac{\partial}{\partial t} \rho_e(\underline{\mathbf{R}}, t)
\end{aligned}$$

2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

We consider the xz plane and assume that the field is independent of y /
Wir betrachten die xz -Ebene und nehmen an, dass das Feld unabhängig von y ist $\Rightarrow \frac{\partial}{\partial y} = 0$

Then it follows for the 3-D wave equations / Es folgt dann für die 3D-Wellengleichungen

$$\begin{aligned} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) &= -\nabla \times \underline{\mathbf{J}}_e(x, z, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(x, z, t) + \frac{1}{\mu_0} \nabla \rho_m(x, z, t) \\ \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) &= \nabla \times \underline{\mathbf{J}}_m(x, z, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(x, z, t) + \frac{1}{\epsilon_0} \nabla \rho_e(x, z, t) \end{aligned}$$

And we confine the current sources to / Und wir beschränken die Stromquellen auf

$$\begin{aligned} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) &= \underline{\mathbf{J}}_m(x, z, t) = J_{my}(x, z, t) \underline{\mathbf{e}}_y \\ \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) &= \underline{\mathbf{J}}_e(x, z, t) = J_{ey}(x, z, t) \underline{\mathbf{e}}_y \end{aligned}$$

This yields for the above given 3-D wave equation /
Dies ergibt für die oben gegebenen 3D-Wellengleichungen

$$\begin{aligned} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) &= \nabla \times [J_{ey}(x, z, t) \underline{\mathbf{e}}_y] + \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \underline{\mathbf{e}}_y + \frac{1}{\mu_0} \nabla \rho_m(x, y, t) \\ \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) &= \nabla \times [J_{my}(x, z, t) \underline{\mathbf{e}}_y] + \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \underline{\mathbf{e}}_y + \frac{1}{\epsilon_0} \nabla \rho_e(x, y, t) \end{aligned}$$

2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

Curl and divergence of the current sources /
Rotation und Divergenz der Stromquellen

$$\begin{aligned} \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) &= J_{my}(x, z, t) \underline{\mathbf{e}}_y \\ \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) &= J_{ey}(x, z, t) \underline{\mathbf{e}}_y \\ \nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) &= \begin{vmatrix} \underline{\mathbf{e}}_x & \underline{\mathbf{e}}_y & \underline{\mathbf{e}}_z \\ \partial_x & \partial_y & \partial_z \\ 0 & J_{ez}(x, z, t) & 0 \end{vmatrix} \\ &= \frac{\partial}{\partial x} J_{ez}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{ez}(x, z, t) \underline{\mathbf{e}}_x \\ &= -\frac{\partial}{\partial z} J_{ez}(x, z, t) \underline{\mathbf{e}}_x + \frac{\partial}{\partial x} J_{ez}(x, z, t) \underline{\mathbf{e}}_z \\ \nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) &= \begin{vmatrix} \underline{\mathbf{e}}_x & \underline{\mathbf{e}}_y & \underline{\mathbf{e}}_z \\ \partial_x & \partial_y & \partial_z \\ 0 & J_{mx}(x, z, t) & 0 \end{vmatrix} \\ &= \frac{\partial}{\partial x} J_{mx}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{mx}(x, z, t) \underline{\mathbf{e}}_x \\ &= -\frac{\partial}{\partial z} J_{mx}(x, z, t) \underline{\mathbf{e}}_x + \frac{\partial}{\partial x} J_{mx}(x, z, t) \underline{\mathbf{e}}_z \end{aligned}$$

$$\begin{aligned} \nabla \cdot \underline{\mathbf{J}}_{my}(x, z, t) \underline{\mathbf{e}}_y &= \left(\underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot J_{my}(x, z, t) \underline{\mathbf{e}}_y = \frac{\partial}{\partial y} J_{my}(x, z, t) = 0 \\ \nabla \cdot \underline{\mathbf{J}}_{ey}(x, z, t) \underline{\mathbf{e}}_y &= \left(\underline{\mathbf{e}}_x \frac{\partial}{\partial x} + \underline{\mathbf{e}}_y \frac{\partial}{\partial y} + \underline{\mathbf{e}}_z \frac{\partial}{\partial z} \right) \cdot J_{ey}(x, z, t) \underline{\mathbf{e}}_y = \frac{\partial}{\partial y} J_{ey}(x, z, t) = 0 \end{aligned}$$

The divergence of the current sources is in this special case zero, because the currents are constant in y direction. / Die Divergenz der Stromquellen ist in diesem speziellen Fall null, da die Ströme in y -Richtung konstant sind.

**2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case /
2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall**

$$\nabla \cdot \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_m(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = -\frac{\partial}{\partial t} \rho_e(\underline{\mathbf{R}}, t)$$

$$\nabla \cdot \underline{\mathbf{J}}_m(x, z, \omega) = j\omega \rho_m(x, z, \omega)$$

$$\nabla \cdot \underline{\mathbf{J}}_e(x, z, \omega) = j\omega \rho_e(x, z, \omega)$$

$$\begin{aligned} \rho_m(x, z, \omega) &= \frac{1}{j\omega} \nabla \cdot \underline{\mathbf{J}}_m(x, z, \omega) \\ \rho_e(x, z, \omega) &= \frac{1}{j\omega} \nabla \cdot \underline{\mathbf{J}}_e(x, z, \omega) \end{aligned}$$

$$\nabla \cdot J_{my}(x, z, \omega) \underline{\mathbf{e}}_y = 0$$

$$\nabla \cdot J_{ey}(x, z, \omega) \underline{\mathbf{e}}_y = 0$$

$$\begin{aligned} \rho_m(\underline{\mathbf{R}}, \omega) &= 0 \\ \rho_e(\underline{\mathbf{R}}, \omega) &= 0 \end{aligned}$$

$$\begin{aligned} \rho_m(\underline{\mathbf{R}}, t) &= 0 \\ \rho_e(\underline{\mathbf{R}}, t) &= 0 \end{aligned}$$

**2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case /
2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall**

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = -\nabla \times \underline{\mathbf{J}}_e(x, z, t) + \epsilon_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_m(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = \nabla \times \underline{\mathbf{J}}_m(x, z, t) + \mu_0 \frac{\partial}{\partial t} \underline{\mathbf{J}}_e(x, z, t)$$

$$\nabla \times \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t) = \frac{\partial}{\partial x} J_{ez}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{ex}(x, z, t) \underline{\mathbf{e}}_x$$

$$\nabla \times \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t) = \frac{\partial}{\partial x} J_{mx}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{mz}(x, z, t) \underline{\mathbf{e}}_x$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{H}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{H}}(x, z, t) = -\frac{\partial}{\partial x} J_{ey}(x, z, t) \underline{\mathbf{e}}_z + \frac{\partial}{\partial z} J_{ey}(x, z, t) \underline{\mathbf{e}}_x + \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \underline{\mathbf{e}}_y$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \underline{\mathbf{E}}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \underline{\mathbf{E}}(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t) \underline{\mathbf{e}}_z - \frac{\partial}{\partial z} J_{my}(x, z, t) \underline{\mathbf{e}}_x + \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \underline{\mathbf{e}}_y$$

2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \mathbf{H}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \mathbf{H}(x, z, t) = - \frac{\partial}{\partial x} J_{ey}(x, z, t) \mathbf{e}_z + \frac{\partial}{\partial z} J_{ey}(x, z, t) \mathbf{e}_x + \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \mathbf{e}_y$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_x(x, z, t) = - \frac{\partial}{\partial z} J_{ey}(x, z, t)$$

Decoupled equations /
Entkoppelte
Gleichungen



$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(x, z, t) = \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_z(x, z, t) = - \frac{\partial}{\partial x} J_{ey}(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) \mathbf{E}(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} \mathbf{E}(x, z, t) = - \frac{\partial}{\partial z} J_{my}(x, z, t) \mathbf{e}_x + \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \mathbf{e}_y + \frac{\partial}{\partial x} J_{my}(x, z, t) \mathbf{e}_z$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(x, z, t) = - \frac{\partial}{\partial z} J_{my}(x, z, t)$$

Decoupled equations /
Entkoppelte
Gleichungen



$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_z(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t)$$

2-D EM Wave Propagation – 2-D TM Case and 2-D TE Case / 2D EM Wellenausbreitung – 2D-TM-Fall und 2D-TE-Fall

→ Separation in 2-D → TM and TE case / TM: transversal magnetic / transversal magnetisch
Separation in 2D → TM- und TE- Fall TE: transversal electric / transversal elektrisch

$$TM_y \text{ case} / TM_y\text{-Fall} \quad \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

$$E_y(x, z, t)$$

$$H_x(x, z, t)$$

$$H_z(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_x(x, z, t) = - \frac{\partial}{\partial z} J_{ey}(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_z(x, z, t) = - \frac{\partial}{\partial x} J_{ey}(x, z, t)$$

$$TE_y \text{ case} / TE_y\text{-Fall} \quad \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(x, z, t) = \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t)$$

$$H_y(x, z, t)$$

$$E_x(x, z, t)$$

$$E_z(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(x, z, t) = - \frac{\partial}{\partial z} J_{my}(x, z, t)$$

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_z(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t)$$

2-D EM Wave Propagation – 2-D TM Case / 2D EM Wellenausbreitung – 2D-TM-Fall



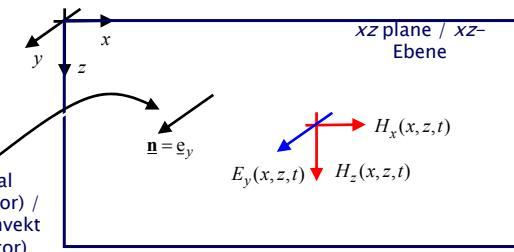
Separation in 2-D → TM case /
Separation in 2D → TM-Fall

TM: transversal magnetic / transversal magnetisch

TM_y case / TM_y-Fall

$$\begin{aligned} E_y(x, z, t) & \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \\ H_x(x, z, t) & \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_x(x, z, t) = - \frac{\partial}{\partial z} J_{ey}(x, z, t) \\ H_z(x, z, t) & \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_z(x, z, t) = - \frac{\partial}{\partial x} J_{ey}(x, z, t) \end{aligned}$$

Surface normal vector (unit-vector) /
Flächennormalenvektor or (Einheitsvektor)



2-D EM Wave Propagation – 2-D TE Case / 2D EM Wellenausbreitung – 2D-TE-Fall



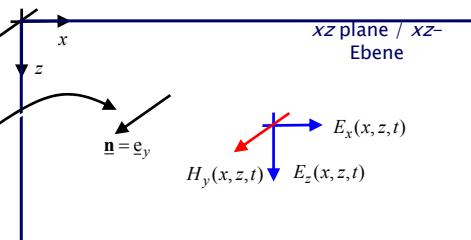
Separation in 2-D → TE case /
Separation in 2D → TE-Fall

TE: transversal electric / transversal elektrisch

TE_y case / TE_y-Fall

$$\begin{aligned} H_y(x, z, t) & \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) H_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} H_y(x, z, t) = \epsilon_0 \frac{\partial}{\partial t} J_{my}(x, z, t) \\ E_x(x, z, t) & \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_x(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_x(x, z, t) = - \frac{\partial}{\partial z} J_{my}(x, z, t) \\ E_z(x, z, t) & \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_z(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_z(x, z, t) = \frac{\partial}{\partial x} J_{my}(x, z, t) \end{aligned}$$

Surface normal vector (unit-vector) /
Flächennormalenvektor or (Einheitsvektor)



FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

Central FD Operators / Zentrale FD-Operatoren

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

Central FD Operators / Zentrale FD-Operatoren

$$\begin{aligned} \frac{\partial^2}{\partial x^2} E_y(x, z, t) &= \frac{E_y(x + \Delta x, z, t) - 2E_y(x, z, t) + E_y(x - \Delta x, z, t)}{(\Delta x)^2} + O[(\Delta x)^2] \\ \frac{\partial^2}{\partial z^2} E_y(x, z, t) &= \frac{E_y(x, z + \Delta z, t) - 2E_y(x, z, t) + E_y(x, z - \Delta z, t)}{(\Delta z)^2} + O[(\Delta z)^2] \\ \frac{\partial^2}{\partial t^2} E_y(x, z, t) &= \frac{E_y(x, z, t + \Delta t) - 2E_y(x, z, t) + E_y(x, z, t - \Delta t)}{(\Delta t)^2} + O[(\Delta t)^2] \end{aligned}$$

Backward FD Operator / Rückwärts-FD-Operator

$$\frac{\partial}{\partial t} J_{ey}(x, z, t) = \frac{J_{ey}(x, z, t) - J_{ey}(x, z, t - \Delta t)}{\Delta t} + O(\Delta t)$$

FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

2-D TM wave equation / 2D-TM-Wellengleichung

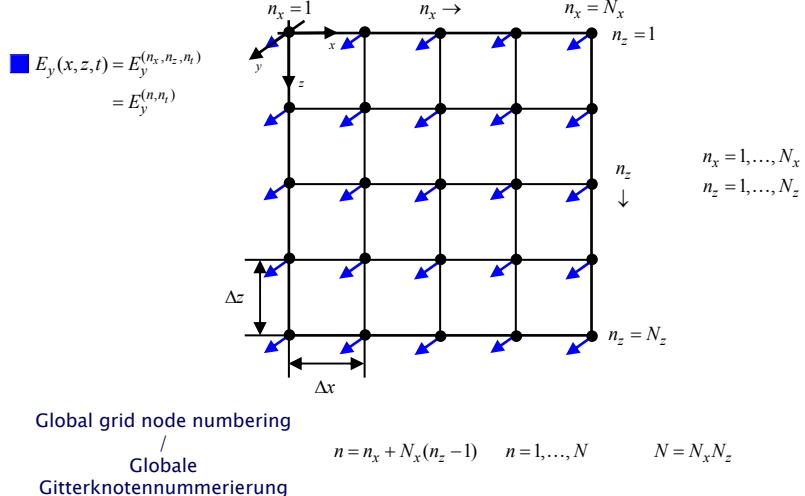
$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t)$$

Explicit FD algorithm in the time domain of 2nd order in space and time / Expliziter FD-Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit

$$\begin{aligned} & \frac{E_y(x + \Delta x, z, t) - 2E_y(x, z, t) + E_y(x - \Delta x, z, t)}{(\Delta x)^2} + \frac{E_y(x, z + \Delta z, t) - 2E_y(x, z, t) + E_y(x, z - \Delta z, t)}{(\Delta z)^2} \\ & - \frac{1}{c_0^2} \frac{E_y(x, z, t + \Delta t) - 2E_y(x, z, t) + E_y(x, z, t - \Delta t)}{(\Delta t)^2} \\ & = \mu_0 \frac{J_{ey}(x, z, t) - J_{ey}(x, z, t - \Delta t)}{\Delta t} + O[(\Delta x)^2] + O[(\Delta z)^2] + O[(\Delta t)^2] \end{aligned}$$

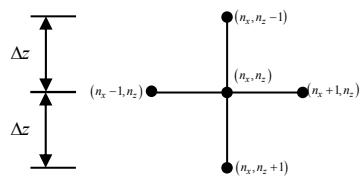
FD Method – 2-D TM Wave Equation – 2-D FD Grid / FD-Methode – 2D-TM-Wellengleichung – 2D-FD-Gitter

2-D FD grid /
2D-FD-Gitter



FD Method – 2-D TM Wave Equation – 2-D FD Stencil / FD-Methode – 2D-TM-Wellengleichung – 2D-FD-Schablone

2-D FD stencil in space /
2D-FD-Schablone im Raum



FD Method – 2–D TM Wave Equation / FD–Methode – 2D–TM–Wellengleichung

**Explicit 2–D FD algorithm in the time domain of 2nd order in space and time /
Expliziter 2D–FD–Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit**

$$\begin{aligned}
 E_y(x, z, t + \Delta t) = & 2E_y(x, z, t) - E_y(x, z, t - \Delta t) \\
 & + c_0^2 \frac{(\Delta t)^2}{(\Delta x)^2} [E_y(x + \Delta x, z, t) - 2E_y(x, z, t) + E_y(x - \Delta x, z, t)] \\
 & + c_0^2 \frac{(\Delta t)^2}{(\Delta z)^2} [E_y(x, z + \Delta z, t) - 2E_y(x, z, t) + E_y(x, z - \Delta z, t)] \\
 & + c_0^2 \mu_0 \Delta t [J_{ey}(x, z, t) - J_{ey}(x, z, t - \Delta t)] + O[(\Delta x)^2] + O[(\Delta z)^2] + O[(\Delta t)^2]
 \end{aligned}$$

Marching-on-in-time algorithm / „Marschieren in der Zeit“–Algorithmus

$$x \rightarrow n_x \Delta x, \quad n_x = 1, \dots, N_x$$

$$z \rightarrow n_z \Delta z, \quad n_z = 1, \dots, N_z$$

$$t \rightarrow n_t \Delta t, \quad n_t = 1, \dots, N_t$$

$$E_y(x, z, t) \rightarrow E_y^{(n_x, n_z, n_t)}$$

$$J_{ey}(x, z, t) \rightarrow J_{ey}^{(n_x, n_z, n_t)}$$

FD Method – 2–D TM Wave Equation / FD–Methode – 2D–TM–Wellengleichung

**Explicit 2–D FD algorithm in the time domain of 2nd order in space and time /
Expliziter 2D–FD–Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit**

$$\begin{aligned}
 E_y^{(n_x, n_z, n_t+1)} = & 2E_y^{(n_x, n_z, n_t)} - E_y^{(n_x, n_z, n_t-1)} \\
 & + \left(\frac{c_0 \Delta t}{\Delta x} \right)^2 [E_y^{(n_x+1, n_z, n_t)} - 2E_y^{(n_x, n_z, n_t)} + E_y^{(n_x-1, n_z, n_t)}] \\
 & + \left(\frac{c_0 \Delta t}{\Delta z} \right)^2 [E_y^{(n_x, n_z+1, n_t)} - 2E_y^{(n_x, n_z, n_t)} + E_y^{(n_x, n_z-1, n_t)}] \\
 & + c_0^2 \mu_0 \Delta t [J_{ey}^{(n_x, n_z, n_t)} - J_{ey}^{(n_x, n_z, n_t-1)}] + O[(\Delta x)^2] + O[(\Delta z)^2] + O[(\Delta t)^2]
 \end{aligned}$$

**Homogeneous 2–D FD grid of quadratic cells /
Homogenes 2D–FD–Gitter aus quadratischen Zellen** $\Delta x = \Delta z$

$$\begin{aligned}
 E_y^{(n_x, n_z, n_t+1)} = & 2E_y^{(n_x, n_z, n_t)} - E_y^{(n_x, n_z, n_t-1)} \\
 & + \left(\frac{c_0 \Delta t}{\Delta x} \right)^2 [E_y^{(n_x, n_z+1, n_t)} + E_y^{(n_x+1, n_z, n_t)} - 4E_y^{(n_x, n_z, n_t)} + E_y^{(n_x-1, n_z, n_t)} + E_y^{(n_x, n_z-1, n_t)}] \\
 & + c_0^2 \mu_0 \Delta t [J_{ey}^{(n_x, n_z, n_t)} - J_{ey}^{(n_x, n_z, n_t-1)}] + O[(\Delta x)^2] + O[(\Delta t)^2]
 \end{aligned}$$

FD Method – 2-D TM Wave Equation / FD-Methode – 2D-TM-Wellengleichung

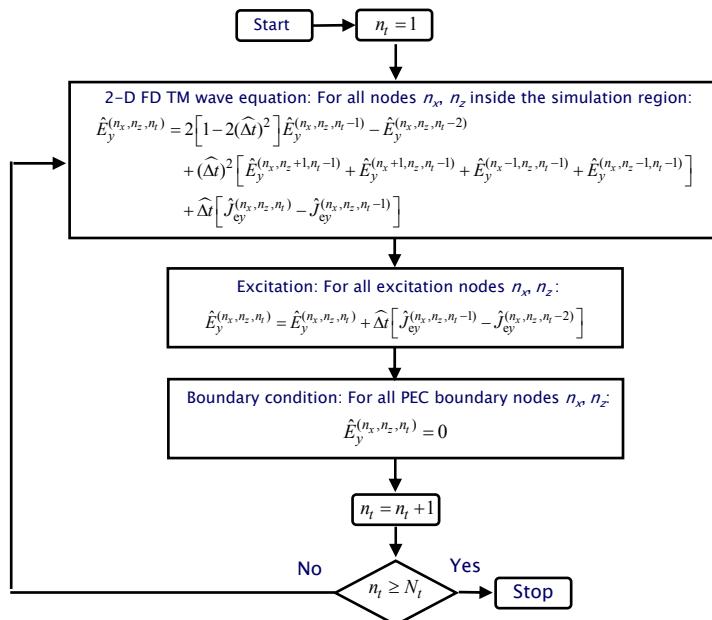
**Explicit FD algorithm in the time domain of 2nd order in space and time /
Expliziter FD-Algorithmus im Zeitbereich 2ter Ordnung in Raum und Zeit**

$$\begin{aligned}\hat{E}_y^{(n_x, n_z, n_t+1)} &= 2\hat{E}_y^{(n_x, n_z, n_t)} - \hat{E}_y^{(n_x, n_z, n_t-1)} \\ &\quad + (\widehat{\Delta t})^2 \left[\hat{E}_y^{(n_x, n_z+1, n_t)} + \hat{E}_y^{(n_x+1, n_z, n_t)} - 4\hat{E}_y^{(n_x, n_z, n_t)} + \hat{E}_y^{(n_x-1, n_z, n_t)} + \hat{E}_y^{(n_x, n_z-1, n_t)} \right] \\ &\quad + \widehat{\Delta t} \left[\hat{J}_{\text{ey}}^{(n_x, n_z, n_t)} - \hat{J}_{\text{ey}}^{(n_x, n_z, n_t-1)} \right]\end{aligned}$$

for / für $\begin{cases} 1 \leq n_x \leq N_x \\ 1 \leq n_z \leq N_z \\ 1 \leq n_t \leq N_t \end{cases}$

$$\begin{aligned}\hat{E}_y^{(n_x, n_z, n_t)} &= 2 \left[1 - 2(\widehat{\Delta t})^2 \right] \hat{E}_y^{(n_x, n_z, n_t-1)} - \hat{E}_y^{(n_x, n_z, n_t-2)} \\ &\quad + (\widehat{\Delta t})^2 \left[\hat{E}_y^{(n_x, n_z+1, n_t)} + \hat{E}_y^{(n_x+1, n_z, n_t)} + \hat{E}_y^{(n_x-1, n_z, n_t)} + \hat{E}_y^{(n_x, n_z-1, n_t)} \right] \\ &\quad + \widehat{\Delta t} \left[\hat{J}_{\text{ey}}^{(n_x, n_z, n_t)} - \hat{J}_{\text{ey}}^{(n_x, n_z, n_t-1)} \right]\end{aligned}$$

FD Method – 2-D FD Wave Equation – TM Case – Flow Chart / FD-Methode – 2D FD-Wellengleichung – TM-Fall – Flussdiagramm



FD Method – 2–D TM Wave Equation – Example / FD-Methode – 2D-TM-Wellengleichung – Beispiel

Scalar 2–D TM wave equation / Skalare 2D-TM-Wellengleichung

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial z^2} \right) E_y(x, z, t) - \frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} E_y(x, z, t) = \mu_0 \frac{\partial}{\partial t} J_{ey}(x, z, t) \quad \text{für } \begin{cases} 0 \leq x \leq X \\ 0 \leq z \leq Z \\ 0 \leq t \leq T \end{cases}$$

Initial condition / Anfangsbedingung

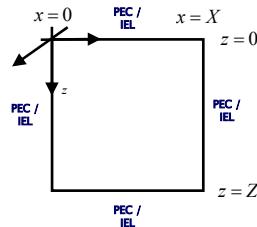
$$E_y(x, z, t) = J_{ey}(x, z, t) = 0 \quad t \leq 0 \quad \text{Causality / Kausalität}$$

$$J_{ey}(x, z, t) = \delta(x - x_0) \delta(z - z_0) f(t) \quad t > 0$$

Boundary conditions for a perfectly electrically conducting (PEC) boundary /
Randbedingung für einen ideal elektrisch leitenden (IEL) Rand

$$\left. \begin{array}{l} E_y(0, z, t) = 0 \\ E_y(X, z, t) = 0 \end{array} \right\} \forall z, t \forall t \quad \text{and / und} \quad \left. \begin{array}{l} E_y(x, 0, t) = 0 \\ E_y(x, Z, t) = 0 \end{array} \right\} \forall x, t \forall t$$

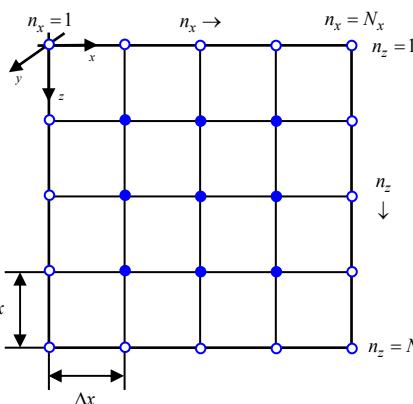
Hyperbolic initial-boundary-value problem /
Hyperbolisches Anfangs-Randwert-Problem



FD Method – 2–D TM Wave Equation – Example / FD-Methode – 2D-TM-Wellengleichung – Beispiel

Nodes in the simulation region / Knoten im Simulationsgebiet

- $E_y^{(n_x, n_z)} \neq 0$



Nodes at the PEC boundary /

Knoten auf dem IEL-Rand

- $E_y^{(n_x, n_z)} = 0$

Global grid node numbering /

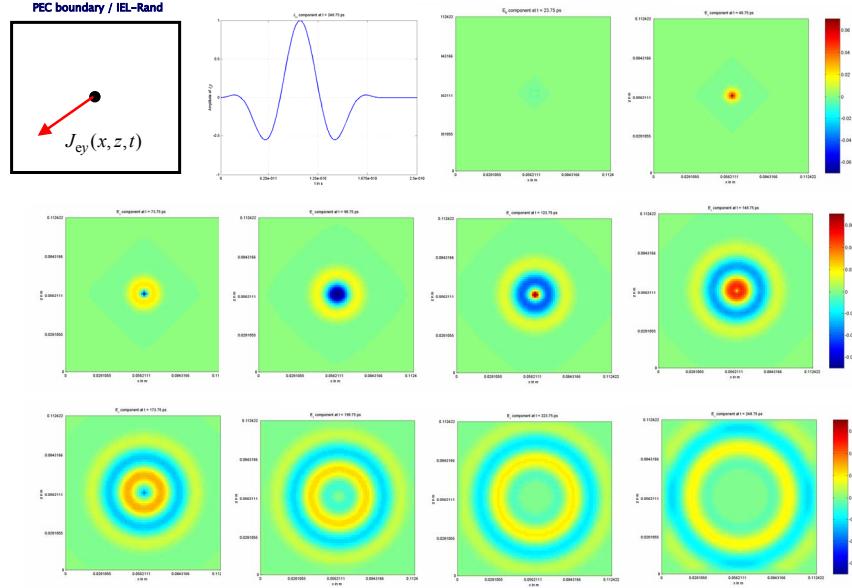
Globale

Gitterknotennummerierung

$$n = n_x + N_x(n_z - 1) \quad n = 1, \dots, N$$

$$N = N_x N_z$$

FD Method – 2-D TM Wave Equation – Example / FD-Methode – 2D-TM-Wellengleichung – Beispiel



FD Method – 2-D FD Wave Equation – TM Case – Validation / FD-Methode – 2D FD-Wellengleichung – TM-Fall – Validierung

Domain integral representation /
(Gebiets-) Integraldarstellung

1-D case / 1D-Fall

$$E_x(z, \omega) = j\omega\mu_0 \int_{z'=-\infty}^{\infty} G(z-z', \omega) J_{\text{ex}}(z', \omega) dz' \quad E_y(x, z, \omega) = j\omega\mu_0 \int_{z'=-\infty}^{\infty} \int_{x'=-\infty}^{\infty} G(x-x', z-z', \omega) J_{\text{ey}}(x', z', \omega) dx' dz'$$

2-D case / 2D-Fall

Green's function / Greensche Funktion

$$G(z, \omega) = \frac{c_0}{2} \left[j \operatorname{PV} \frac{1}{\omega_0} + \pi \delta(z) \right] e^{jk_0|z|}$$

$$G(z, t) = \frac{c_0}{2} u \left(t - \frac{|z|}{c_0} \right)$$

$$G(x, z, \omega) = \frac{j}{4} H_0^{(1)} \left(\frac{\omega}{c_0} \sqrt{x^2 + z^2} \right)$$

$$G(x, z, t) = \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - (x^2 + z^2)}} u \left(t - \frac{\sqrt{x^2 + z^2}}{c_0} \right)$$

FD Method – 2-D FD Wave Equation – TM Case – Validation /
FD-Methode – 2D FD-Wellengleichung – TM-Fall – Validierung

2-D Domain integral representation /
 2D-(Gebiets-)Integraldarstellung

$$E_y(r, \omega) = j\omega \mu_0 \int_{r'=0}^{\infty} G(r-r', \omega) J_{ey}(r', \omega) dr'$$

$$G(r, \omega) = \frac{j}{4} H_0^{(1)} \left(\frac{\omega}{c_0} r \right)$$

$$G(r, t) = \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - r^2}} u \left(t - \frac{r}{c_0} \right)$$

$$E_y(\underline{r}, \omega) = j\omega \mu_0 \int_{\underline{r}'} G(\underline{r} - \underline{r}', \omega) J_{ey}(\underline{r}', \omega) d\underline{r}'$$

$$G(\underline{r} - \underline{r}', \omega) = \frac{j}{4} H_0^{(1)} \left(\frac{\omega}{c_0} |\underline{r} - \underline{r}'| \right)$$

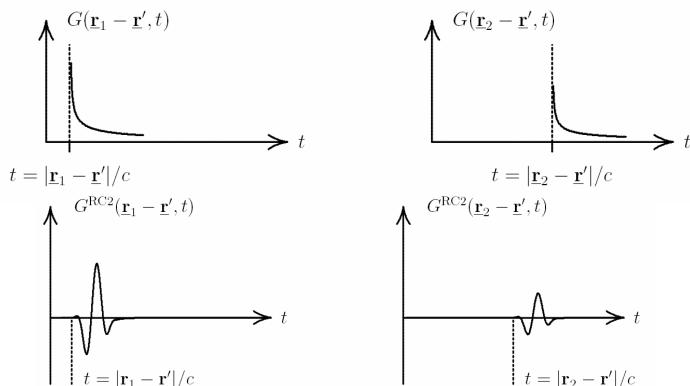
$$G(\underline{r} - \underline{r}', t) = \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - |\underline{r} - \underline{r}'|^2}} u \left(t - \frac{|\underline{r} - \underline{r}'|}{c_0} \right)$$

FD Method – 2-D FD Wave Equation – TM Case – Validation /
FD-Methode – 2D FD-Wellengleichung – TM-Fall – Validierung

2-D Domain integral representation / 2D-(Gebiets-)Integraldarstellung

$$G^{\text{RC2}}(\underline{r} - \underline{r}', \omega) = \text{RC2}(\omega) \frac{j}{4} H_0^{(1)} \left(\frac{\omega}{c_0} |\underline{r} - \underline{r}'| \right)$$

$$G^{\text{RC2}}(\underline{r} - \underline{r}', t) = \text{RC2}(t) *_t \frac{c_0}{2\pi} \frac{1}{\sqrt{c_0^2 t^2 - |\underline{r} - \underline{r}'|^2}} u \left(t - \frac{|\underline{r} - \underline{r}'|}{c_0} \right)$$



$$G^{\text{RC2}}(\underline{r} - \underline{r}', \omega) \approx \frac{1}{4} e^{\frac{j\pi}{4}} \sqrt{\frac{2c}{\pi}} \frac{\text{RC2}(\omega)}{\sqrt{\omega}} \frac{e^{j k |\underline{r} - \underline{r}'|}}{\sqrt{|\underline{r} - \underline{r}'|}}$$

EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

The first two Maxwell's Equations are: /
Die ersten beiden Maxwellschen Gleichungen lauten:

Equations of first order /
Gleichungen der ersten Ordnung

$$\frac{\partial}{\partial t} \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

Constitutive Equations for Vacuum /
Konstituierende Gleichungen
(Materialgleichungen) für Vakuum

$$\underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \mu \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \epsilon \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

$$f(\underline{\mathbf{H}}, \underline{\mathbf{E}})$$

Constitutive Equations for Vacuum /
Konstituierende Gleichungen
(Materialgleichungen) für Vakuum

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = \nu_0 \underline{\mathbf{B}}(\underline{\mathbf{R}}, t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

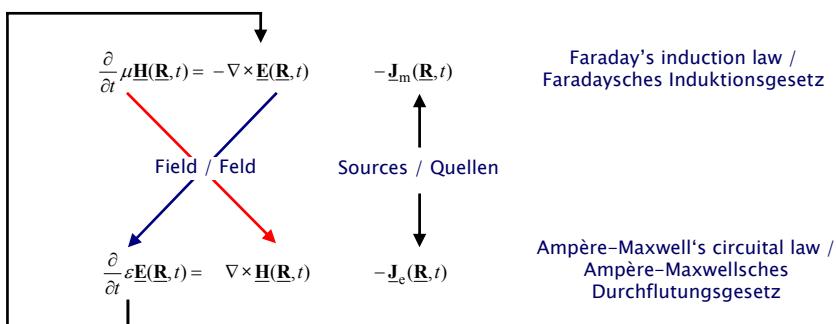
$$\frac{\partial}{\partial t} \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} [\epsilon \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)] = \nabla \times [\nu \underline{\mathbf{B}}(\underline{\mathbf{R}}, t)] - \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

$$f(\underline{\mathbf{B}}, \underline{\mathbf{E}})$$

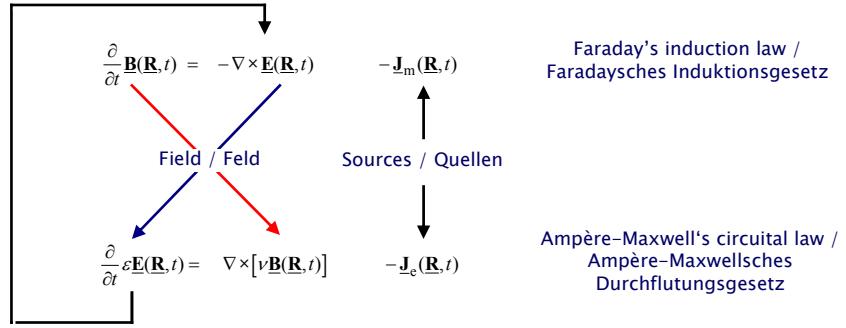
EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

Idea: Outline of a flow chart /
Idee: Entwurf eines Flussdiagramms



EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

Idea: Outline of a flow chart /
Idee: Entwurf eines Flussdiagramms



1-D EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / 1D EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

The first two Maxwell's Equations are: /
Die ersten beiden Maxwellschen Gleichungen lauten:

Constitutive Equations for Vacuum /
Konstituierende Gleichungen
(Materialgleichungen) für Vakuum

$$\underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = \mu_0 \underline{\mathbf{H}}(\underline{\mathbf{R}}, t)$$

$$\underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \epsilon_0 \underline{\mathbf{E}}(\underline{\mathbf{R}}, t)$$

Ansatz for the electric and
magnetic field strength /
Ansatz für die elektrische und
magnetische Feldstärke

$$\underline{\mathbf{E}}(\underline{\mathbf{R}}, t) = E_x(z, t) \mathbf{e}_x$$

$$\underline{\mathbf{H}}(\underline{\mathbf{R}}, t) = H_y(z, t) \mathbf{e}_y$$

$$\frac{\partial}{\partial t} \underline{\mathbf{B}}(\underline{\mathbf{R}}, t) = -\nabla \times \underline{\mathbf{E}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_m(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} \underline{\mathbf{D}}(\underline{\mathbf{R}}, t) = \nabla \times \underline{\mathbf{H}}(\underline{\mathbf{R}}, t) - \underline{\mathbf{J}}_e(\underline{\mathbf{R}}, t)$$

$$\frac{\partial}{\partial t} H_y(z, t) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z, t) - \frac{1}{\mu_0} J_{my}(z, t)$$

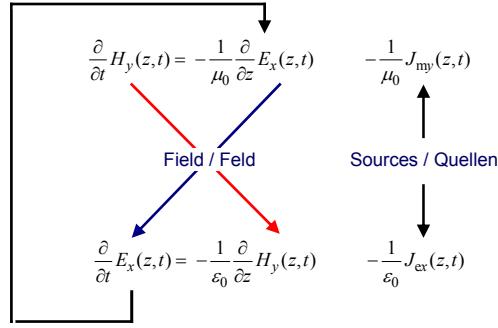
$$\frac{\partial}{\partial t} E_x(z, t) = -\frac{1}{\epsilon_0} \frac{\partial}{\partial z} H_y(z, t) - \frac{1}{\epsilon_0} J_{ex}(z, t)$$

$$\frac{d}{dt} f(t) = \frac{f\left(t + \frac{\Delta t}{2}\right) - f\left(t - \frac{\Delta t}{2}\right)}{\Delta t} + O[(\Delta t)^2]$$

$$\frac{d}{dz} f(z) = \frac{f\left(z + \frac{\Delta z}{2}\right) - f\left(z - \frac{\Delta z}{2}\right)}{\Delta x} + O[(\Delta x)^2]$$

1-D EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / 1D EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

Idea: Outline of a flow chart /
Idee: Entwurf eines Flussdiagramms



1-D EM Wave Propagation – FDTD – Discretization of the 1st Equation / 1D EM Wellenausbreitung – FDTD – Diskretisierung der 1ten Gleichung

Spatial discretization of the 1st equation /
Räumliche Diskretisierung der 1ten Gleichung

$$\frac{\partial}{\partial t} H_y(z,t) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z,t) - \frac{1}{\mu_0} J_{my}(z,t)$$

$$H_y : z \rightarrow n_z \Delta z, \quad n_z = 1, \dots, N_z$$

$$E_x : z \rightarrow (n_z + 1/2) \Delta z, \quad n_z = 1, \dots, N_z$$

$$\frac{\partial}{\partial z} E_x(z,t) \rightarrow \left. \frac{\partial}{\partial z} E_x(z,t) \right|_z = \frac{1}{\Delta z} \left[E_x \left(z + \frac{\Delta z}{2} \right) - E_x \left(z - \frac{\Delta z}{2} \right) \right] + O[(\Delta z)^2]$$

$$\begin{bmatrix} E_x^{(n_z)} \\ E_x^{(n_z-1/2)} & H_y^{(n_z)} & E_x^{(n_z+1/2)} \end{bmatrix}$$

$$n_z - \frac{1}{2}, \quad n_z, \quad n_z + \frac{1}{2}$$

$$\frac{\partial}{\partial t} H_y^{(n_z)}(t) = -\frac{1}{\mu_0} \frac{1}{\Delta z} [E_x^{(n_z+1/2)}(t) - E_x^{(n_z-1/2)}(t)] - \frac{1}{\mu_0} J_{my}^{(n_z)}(t)$$

1-D EM Wave Propagation – FDTD – Discretization of the 2nd Equation / 1D EM Wellenausbreitung – FDTD – Diskretisierung der 2ten Gleichung

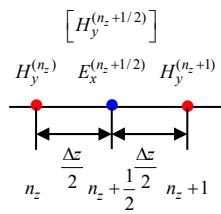
Spatial discretization of the 2nd equation /
Räumliche Diskretisierung der 2ten Gleichung

$$\frac{\partial}{\partial t} E_x(z, t) = -\frac{1}{\epsilon_0} \frac{\partial}{\partial z} H_y(z, t) - \frac{1}{\epsilon_0} J_{ex}(z, t)$$

$$H_y : z \rightarrow n_z \Delta z, \quad n_z = 1, \dots, N_z$$

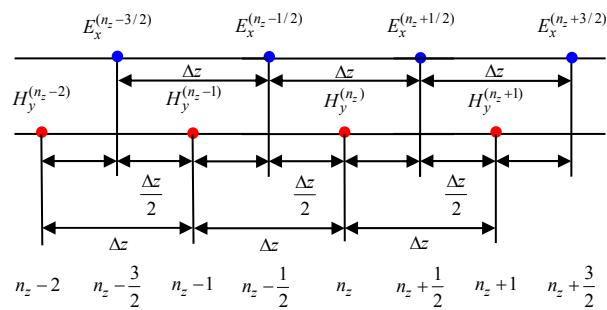
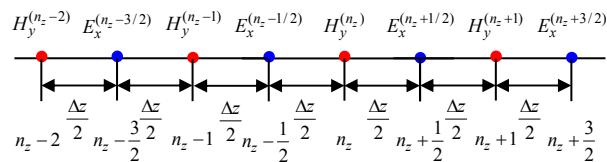
$$E_x : z \rightarrow (n_z + 1/2) \Delta z, \quad n_z = 1, \dots, N_z$$

$$\frac{\partial}{\partial z} H_y(z, t) \rightarrow \left. \frac{\partial}{\partial z} H_y(z, t) \right|_{z+\frac{\Delta z}{2}} = \frac{1}{\Delta z} [H_y(z + \Delta z) - H_y(z)] + O((\Delta z)^2)$$



$$\frac{\partial}{\partial t} E_x^{(n_z+1/2)}(t) = -\frac{1}{\epsilon_0 \Delta z} [H_y^{(n_z+1)}(t) - H_y^{(n_z)}(t)] - \frac{1}{\epsilon_0} J_{ex}^{(n_z+1/2)}(t)$$

1-D EM Wave Propagation – 1-D FDTD – Staggered Grid in Space / 1D EM Wellenausbreitung – 1-D FDTD – Versetztes Gitter im Raum



1-D EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / 1D EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

$$\frac{\partial}{\partial t} H_y(z, t) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z, t) - \frac{1}{\mu_0} J_{my}(z, t)$$

$$\frac{\partial}{\partial t} E_x(z, t) = -\frac{1}{\epsilon_0} \frac{\partial}{\partial z} H_y(z, t) - \frac{1}{\epsilon_0} J_{ex}(z, t)$$



$$\frac{d}{dz} f(z) = \frac{1}{\Delta z} \left[f\left(z + \frac{\Delta z}{2}\right) - f\left(z - \frac{\Delta z}{2}\right) \right] + O[(\Delta z)^2]$$



$$\frac{\partial}{\partial t} H_y^{(n_z)}(t) = -\frac{1}{\mu_0} \frac{1}{\Delta z} \left[E_x^{(n_z+1/2)}(t) - E_x^{(n_z-1/2)}(t) \right] - \frac{1}{\mu_0} J_{my}^{(n_z)}(t)$$

$$\frac{\partial}{\partial t} E_x^{(n_z+1/2)}(t) = -\frac{1}{\epsilon_0 \Delta z} \left[H_y^{(n_z+1)}(t) - H_y^{(n_z)}(t) \right] - \frac{1}{\epsilon_0} J_{ex}^{(n_z+1/2)}(t)$$

$$\frac{\partial}{\partial t} H_y^{(n_z)}(t) = ?$$

$$\frac{\partial}{\partial t} E_x^{(n_z+1/2)}(t) = ?$$

1-D EM Wave Propagation – Finite-Difference Time-Domain (FDTD) / 1D EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

$$\frac{\partial}{\partial t} H_y^{(n_z)}(t) = -\frac{1}{\mu_0} \frac{1}{\Delta z} \left[E_x^{(n_z+1/2)}(t) - E_x^{(n_z-1/2)}(t) \right] - \frac{1}{\mu_0} J_{my}^{(n_z)}(t)$$

$$\frac{\partial}{\partial t} E_x^{(n_z+1/2)}(t) = -\frac{1}{\epsilon_0 \Delta z} \left[H_y^{(n_z+1)}(t) - H_y^{(n_z)}(t) \right] - \frac{1}{\epsilon_0} J_{ey}^{(n_z+1/2)}(t)$$

$$\frac{d}{dt} f(t) = \frac{1}{\Delta t} \left[f\left(t + \frac{\Delta t}{2}\right) - f\left(t - \frac{\Delta t}{2}\right) \right] + O[(\Delta t)^2]$$

Staggered grid in time / Versetztes Gitter in der Zeit

$$\frac{\partial}{\partial t} H_y^{(n_z)}(t) = \frac{H_y^{(n_z, n_t)} - H_y^{(n_z, n_t-1)}}{\Delta t} + O[(\Delta t)^2]$$

$$\frac{\partial}{\partial t} E_x^{(n_z+1/2)}(t) = \frac{E_x^{(n_z+1/2, n_t+1/2)} - E_x^{(n_z+1/2, n_t-1/2)}}{\Delta t} + O[(\Delta t)^2]$$

$$\frac{H_y^{(n_z, n_t)} - H_y^{(n_z, n_t-1)}}{\Delta t} = -\frac{1}{\mu_0} \frac{1}{\Delta z} \left[E_x^{(n_z+1/2)}(t) - E_x^{(n_z-1/2)}(t) \right] - \frac{1}{\mu_0} J_{my}^{(n_z)}(t)$$

$$\frac{E_x^{(n_z+1/2, n_t+1/2)} - E_x^{(n_z+1/2, n_t-1/2)}}{\Delta t} = -\frac{1}{\epsilon_0 \Delta z} \left[H_y^{(n_z+1)}(t) - H_y^{(n_z)}(t) \right] - \frac{1}{\epsilon_0} J_{ey}^{(n_z+1/2)}(t)$$

1-D EM Wave Propagation – Finite–Difference Time–Domain (FDTD) / 1D EM Wellenausbreitung – Finite Differenzen im Zeitbereich (FDTD)

$$\frac{H_y^{(n_z, n_t)} - H_y^{(n_z, n_t-1)}}{\Delta t} = -\frac{1}{\mu_0 \Delta z} \left[E_x^{(n_z+1/2)}(t) - E_x^{(n_z-1/2)}(t) \right] - \frac{1}{\mu_0} J_{\text{my}}^{(n_z)}(t)$$

$$\frac{E_x^{(n_z+1/2, n_t+1/2)} - E_x^{(n_z+1/2, n_t+1/2)}}{\Delta t} = -\frac{1}{\epsilon_0 \Delta z} \left[H_y^{(n_z+1)}(t) - H_y^{(n_z)}(t) \right] - \frac{1}{\epsilon_0} J_{\text{ey}}^{(n_z+1/2)}(t)$$

Explicit 1-D FDTD algorithm on a staggered grid in space and time /
Expliziter 1D-FDTD–Algorithmus auf einem versetzten Gitter im Raum und Zeit

$$H_y^{(n_z, n_t)} = H_y^{(n_z, n_t-1)} - \frac{\Delta t}{\mu_0 \Delta z} \left[E_x^{(n_z+1/2, n_t-1/2)} - E_x^{(n_z-1/2, n_t-1/2)} \right] - \frac{\Delta t}{\mu_0} J_{\text{my}}^{(n_z, n_t-1/2)}$$

$$E_x^{(n_z+1/2, n_t+1/2)} = E_x^{(n_z+1/2, n_t+1/2)} - \frac{\Delta t}{\epsilon_0 \Delta z} \left[H_y^{(n_z+1, n_t)} - H_y^{(n_z, n_t)} \right] - \frac{\Delta t}{\epsilon_0} J_{\text{ey}}^{(n_z+1/2, n_t)}$$

FDTD: Yee, K. S.: Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media. *IEEE Transactions on Antennas Propagation*, Vol. AP-14, pp. 302–307, 1966.

1-D EM Wave Propagation – 1-D FDTD / 1D EM Wellenausbreitung – 1D FDTD

The first two Maxwell's Equations are: /
Die ersten beiden Maxwellschen Gleichungen lauten:

$$\frac{\partial}{\partial t} H_y(z, t) = -\frac{1}{\mu_0} \frac{\partial}{\partial z} E_x(z, t) - \frac{1}{\mu_0} J_{\text{my}}(z, t)$$

$$\frac{\partial}{\partial t} E_x(z, t) = -\frac{1}{\epsilon_0} \frac{\partial}{\partial z} H_y(z, t) - \frac{1}{\epsilon_0} J_{\text{ex}}(z, t)$$

Explicit 1-D FDTD algorithm of leap-frog type on a staggered grid in space and time /
Expliziter 1D-FDTD–Algorithmus vom „Bocksprung“–Typ auf einem versetzten Gitter im Raum und Zeit

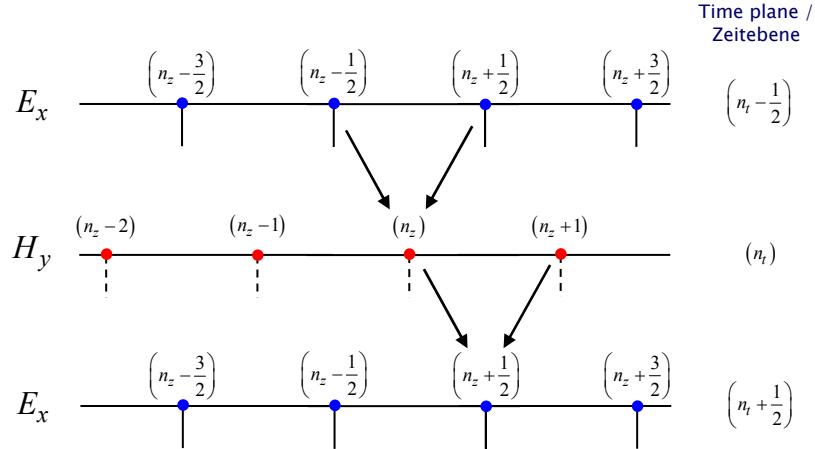
$$H_y^{(n_z, n_t)} = H_y^{(n_z, n_t-1)} - \frac{\Delta t}{\mu_0 \Delta z} \left[E_x^{(n_z+1/2, n_t-1/2)} - E_x^{(n_z-1/2, n_t-1/2)} \right] - \frac{\Delta t}{\mu_0} J_{\text{my}}^{(n_z, n_t-1/2)}$$

$$E_x^{(n_z+1/2, n_t+1/2)} = E_x^{(n_z+1/2, n_t-1/2)} - \frac{\Delta t}{\epsilon_0 \Delta z} \left[H_y^{(n_z+1, n_t)} - H_y^{(n_z, n_t)} \right] - \frac{\Delta t}{\epsilon_0} J_{\text{ex}}^{(n_z+1/2, n_t)}$$

FDTD: Yee, K. S.: Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media. *IEEE Transactions on Antennas Propagation*, Vol. AP-14, pp. 302–307, 1966.

1-D EM Wave Propagation – 1-D FDTD – Staggered Grid in Space / 1D EM Wellenausbreitung – 1-D FDTD – Versetztes Gitter im Raum

Interleaving of the E_x and H_y field components in space and time in the 1-D FDTD formulation /
Überlappung der E_x - und H_y -Feldkomponente in der 1D-FDTD-Formulierung im Raum und in der
Zeit



1-D EM Wave Propagation – FDTD – Normalization / 1D EM Wellenausbreitung – FDTD – Normierung

$$H_y^{(n_z, n_t)} = H_y^{(n_z, n_t-1)} - \frac{\Delta t}{\mu_0 \Delta z} [E_x^{(n_z+1/2, n_t-1/2)} - E_x^{(n_z-1/2, n_t-1/2)}] - \frac{\Delta t}{\mu_0} J_{my}^{(n_z, n_t-1/2)}$$

$$E_x^{(n_z+1/2, n_t+1/2)} = E_x^{(n_z+1/2, n_t-1/2)} - \frac{\Delta t}{\epsilon_0 \Delta z} [H_y^{(n_z+1, n_t)} - H_y^{(n_z, n_t)}] - \frac{\Delta t}{\epsilon_0} J_{ex}^{(n_z+1/2, n_t)}$$

$$\Delta t = \Delta_{\text{ref}} \hat{\Delta t} \quad \Delta t_{\text{ref}} = \frac{\Delta x_{\text{ref}}}{c_{\text{ref}}} \quad \Delta t = \frac{\Delta x_{\text{ref}}}{c_{\text{ref}}} \hat{\Delta t}$$

$$\Delta z = \Delta x_{\text{ref}} \hat{\Delta z} \quad c = c_{\text{ref}} \hat{c} \quad \epsilon = \epsilon_{\text{ref}} \hat{\epsilon} \quad \mu = \mu_{\text{ref}} \hat{\mu} \quad \mu_{\text{ref}} = \mu_0$$

$$E_x = E_{\text{ref}} \hat{E}_x$$

$$H_y = H_{\text{ref}} \hat{H}_y \quad H_{\text{ref}} = \frac{E_{\text{ref}}}{c_{\text{ref}} \mu_{\text{ref}}} = \frac{\sqrt{\epsilon_{\text{ref}} \mu_{\text{ref}}}}{\mu_{\text{ref}}} E_{\text{ref}} = \sqrt{\frac{\epsilon_{\text{ref}}}{\mu_{\text{ref}}}} E_{\text{ref}} = \frac{E_{\text{ref}}}{Z_{\text{ref}}}$$

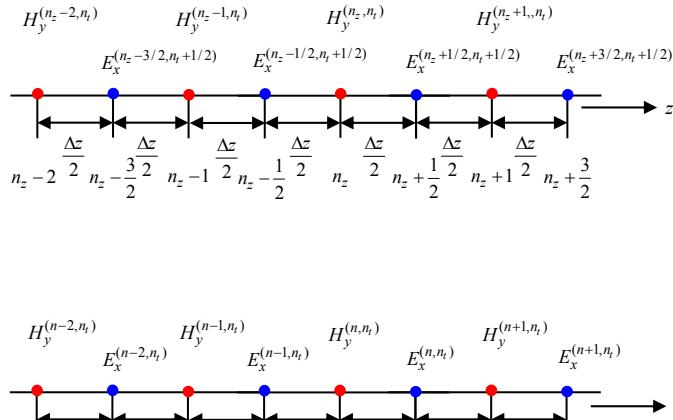
$$J_{\text{ex}} = J_{\text{e ref}} \hat{J}_{\text{ex}} \quad J_{\text{e ref}} = \frac{\epsilon_{\text{ref}}}{\Delta t_{\text{ref}}} E_{\text{ref}}$$

$$J_{\text{mx}} = J_{\text{m ref}} \hat{J}_{\text{mx}} \quad J_{\text{m ref}} = \frac{\mu_{\text{ref}}}{\Delta t_{\text{ref}}} H_{\text{ref}} = \frac{E_{\text{ref}}}{\Delta t_{\text{ref}} c_{\text{ref}}}$$

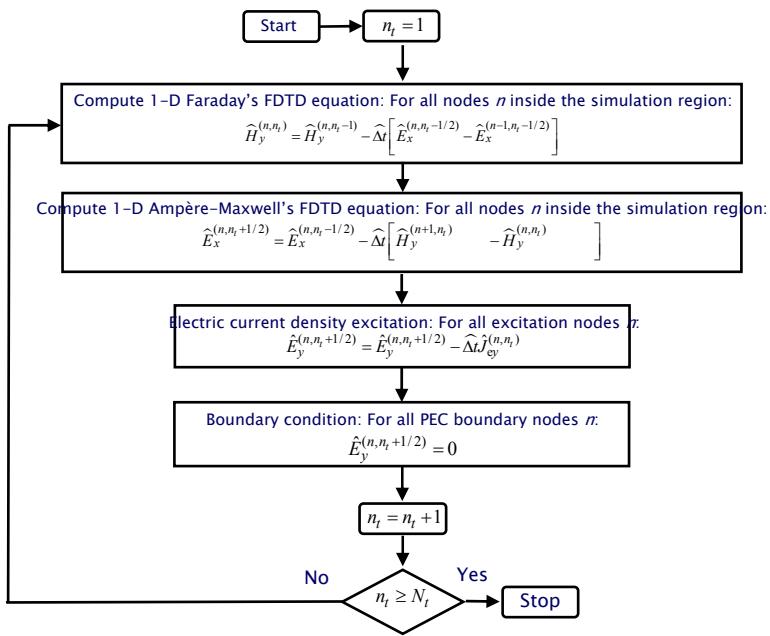
$$\hat{H}_y^{(n_z, n_t)} = \hat{H}_y^{(n_z, n_t-1)} - \hat{\Delta t} [\hat{E}_x^{(n_z+1/2, n_t-1/2)} - \hat{E}_x^{(n_z-1/2, n_t-1/2)}] - \hat{\Delta t} \hat{J}_{\text{my}}^{(n_z, n_t-1/2)}$$

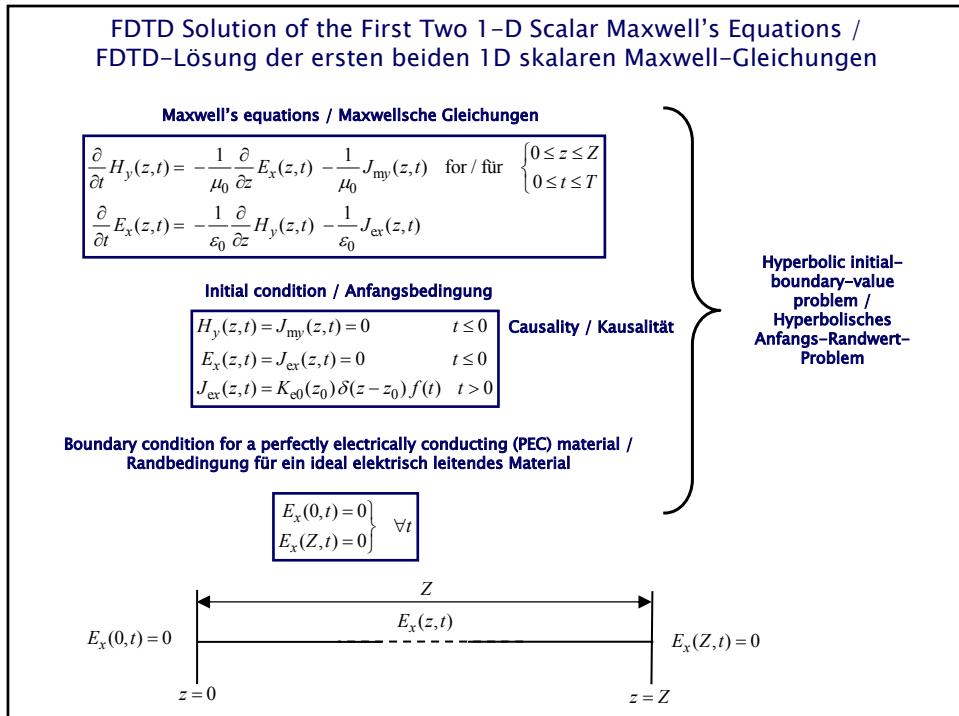
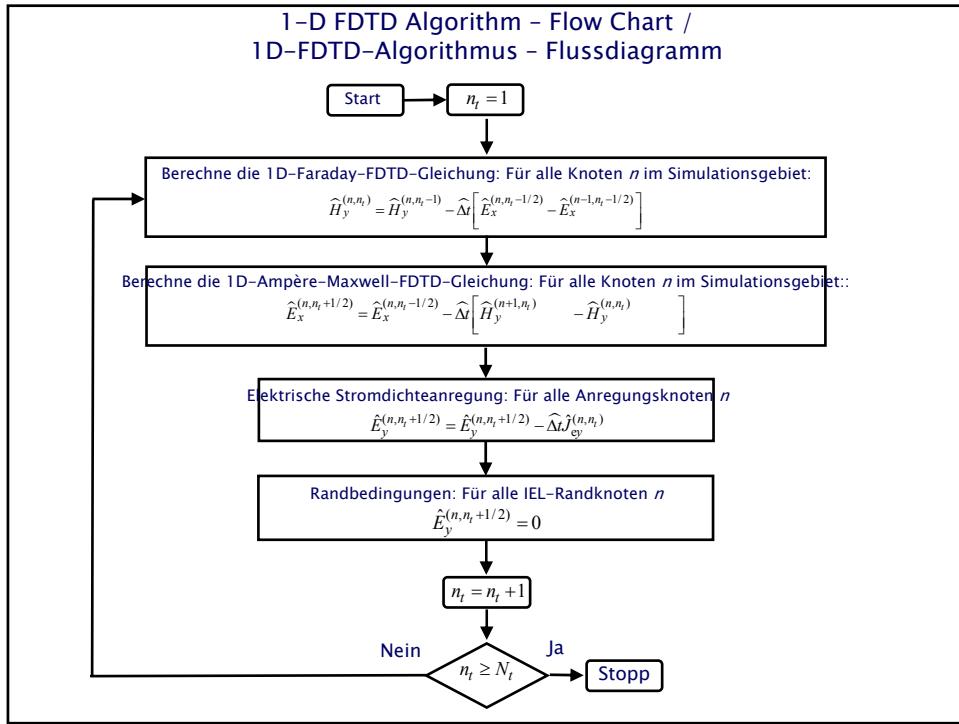
$$\hat{E}_x^{(n_z+1/2, n_t+1/2)} = \hat{E}_x^{(n_z+1/2, n_t-1/2)} - \hat{\Delta t} [\hat{H}_y^{(n_z+1, n_t)} - \hat{H}_y^{(n_z, n_t)}] - \hat{\Delta t} \hat{J}_{\text{ex}}^{(n_z+1/2, n_t)}$$

**1-D FDTD – Staggered Grid in Space – Global Node Numbering /
1D-FDTD – Versetztes Gitter im Raum – Globale Knotennummerierung**



**1-D FDTD Algorithm – Flow Chart /
1D-FDTD-Algorithmus – Flussdiagramm**





FDTD Solution of the First Two 1-D Scalar Maxwell's Equations / FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen

Discrete 1-D FDTD equations / Diskrete 1D-FDTD-Gleichungen

$$\begin{aligned}\hat{H}_y^{(n_z, n_t)} &= \hat{H}_y^{(n_z, n_t-1)} - \Delta t \left[\hat{E}_x^{(n_z+1/2, n_t-1/2)} - \hat{E}_x^{(n_z-1/2, n_t-1/2)} \right] - \Delta t \hat{J}_{\text{my}}^{(n_z, n_t-1/2)} \quad \text{für } \begin{cases} 1 \leq n_z \leq N_z \\ 1 \leq n_t \leq N_t \end{cases} \\ \hat{E}_x^{(n_z+1/2, n_t+1/2)} &= \hat{E}_x^{(n_z+1/2, n_t-1/2)} - \Delta t \left[\hat{H}_y^{(n_z+1, n_t)} - \hat{H}_y^{(n_z, n_t)} \right] - \Delta t \hat{J}_{\text{ex}}^{(n_z+1/2, n_t)}\end{aligned}$$

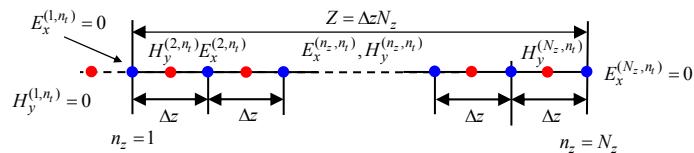
Initial condition / Anfangsbedingung

$$\begin{aligned}H_y^{(n_z, n_t)} &= J_{\text{my}}^{(n_z, n_t)} = 0 & n_t \leq 1 \\ E_x^{(n_z, n_t)} &= J_{\text{ex}}^{(n_z, n_t)} = 0 & n_t \leq 1 \\ J_{\text{ex}}^{(n_z, n_t)} &= K_{\text{ex}}^{(n_z_0)} \delta^{(n_z - n_z_0)} f^{(n_t)} & n_t > 1\end{aligned}$$

Causality / Kausalität

Boundary condition for a perfectly electrically conducting (PEC) material / Randbedingung für ein ideal elektrisch leitendes Material

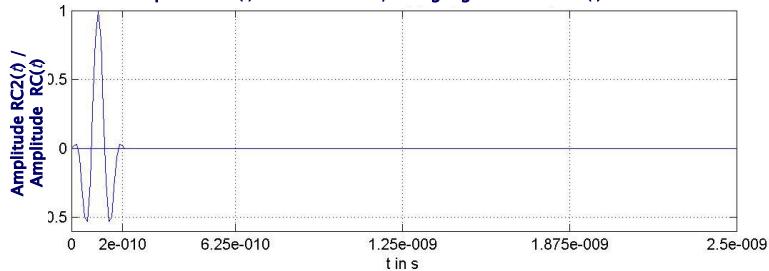
$$\begin{cases} E_x^{(1, n_t)} = 0 \\ E_x^{(N_z, n_t)} = 0 \end{cases} \quad 1 \leq n_t \leq N_t$$



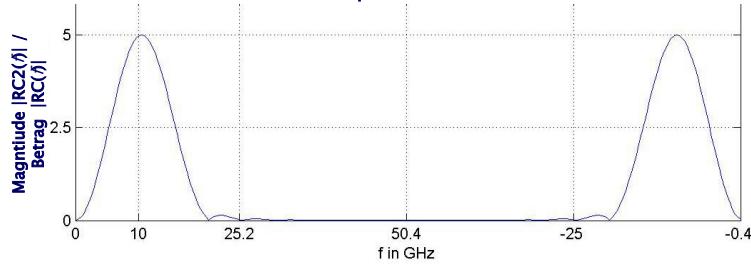
Discrete hyperbolic
initial-boundary-value
problem /
Diskretes
hyperbolisches
Anfangs-Randwert-
Problem

FDTD Solution of the First Two 1-D Scalar Maxwell's Equations / FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen

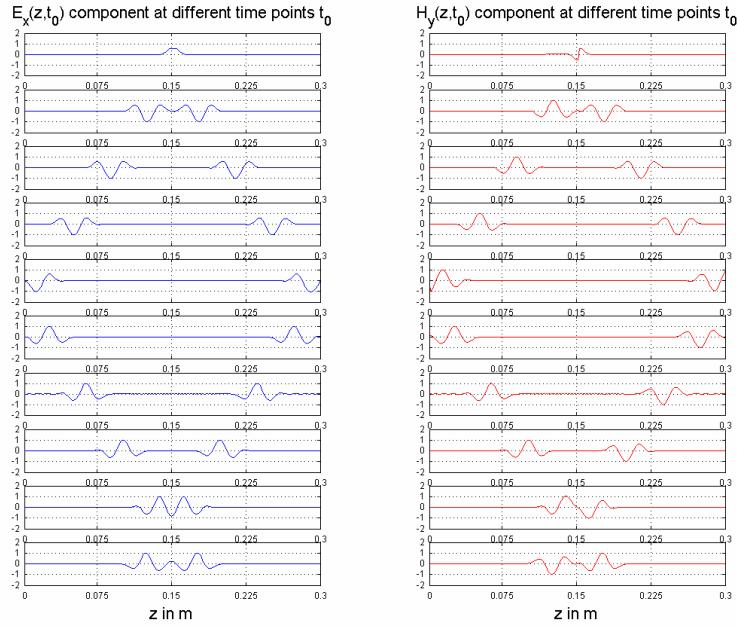
Excitation pulse: RC2(t) – Time Domain / Anregungsfunktion: RC2(t) – Zeitbereich



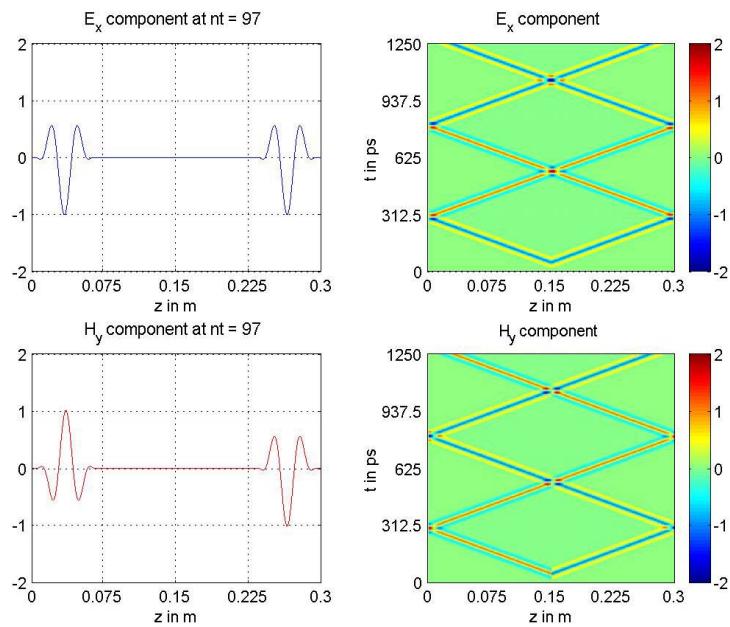
Excitation pulse: RC2(f) – Frequency Domain / Anregungsfunktion: RC(f) – Frequenzbereich



FDTD Solution of the First Two 1-D Scalar Maxwell's Equations /
FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen



FDTD Solution of the First Two 1-D Scalar Maxwell's Equations /
FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen



Implementation of Boundary Conditions / Implementierung von Randbedingungen

**Boundary condition for a perfectly electrically conducting (PEC) material /
Randbedingung für ein ideal elektrisch leitendes Material**

$$\left. \begin{array}{l} E_x^{(1,n_t)} = 0 \\ E_x^{(N_z,n_t)} = 0 \end{array} \right\} \quad 1 \leq n_t \leq N_t$$

**Absorbing/open boundary condition /
Absorbierende/offene Randbedingung**

**Space-time-extrapolation of the first order /
Raum-Zeit-Extrapolation der ersten Ordnung**

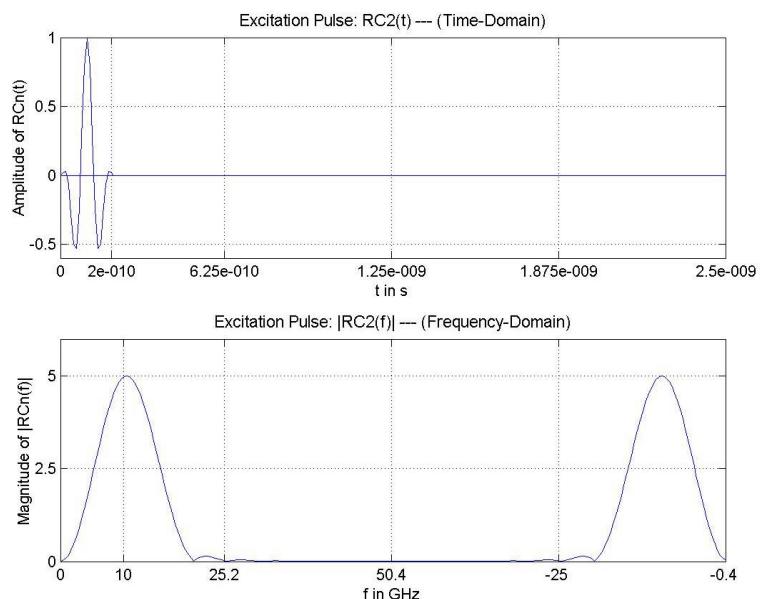
For / Für $\hat{\Delta t} = 0.5$

a plane wave needs two time steps, $2 n_t$, to travel over one grid cell with the size Δz /
braucht eine ebene Welle zwei Zeitschritte, $2 n_t$, um sich über eine Gitterzelle der Größe Δz
auszubreiten

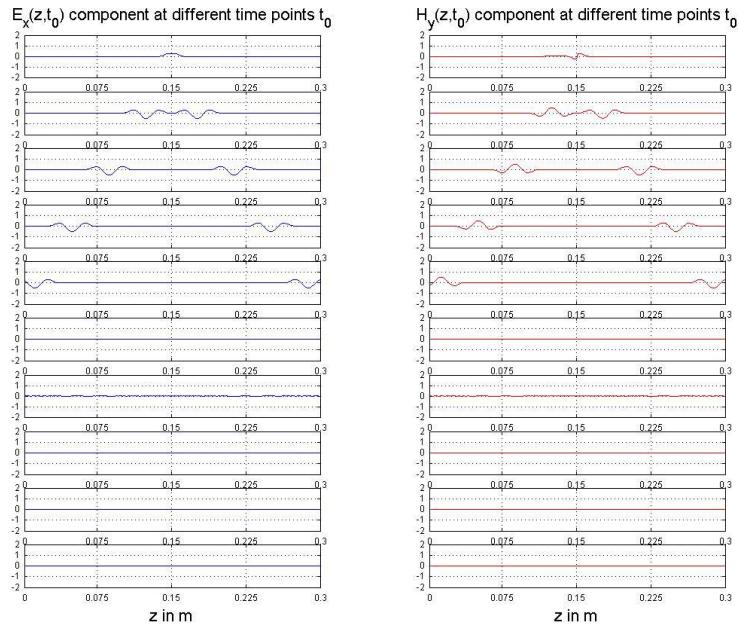
$$\left. \begin{array}{l} E_x^{(1,n_t)} = E_x^{(2,n_t-2)} \\ E_x^{(N_z,n_t)} = E_x^{(N_z-1,n_t-2)} \end{array} \right\} \quad 1 \leq n_t \leq N_t$$

**Space-time-extrapolation of the first order /
Raum-Zeit-Extrapolation der ersten Ordnung**

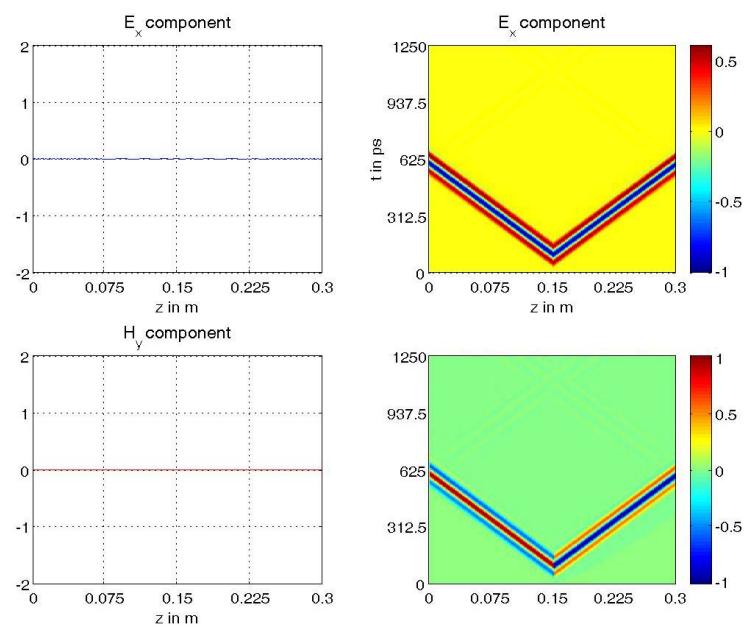
FDTD Solution of the First Two 1-D Scalar Maxwell's Equations / FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen



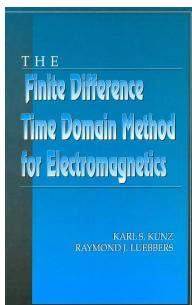
FDTD Solution of the First Two 1-D Scalar Maxwell's Equations / FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen



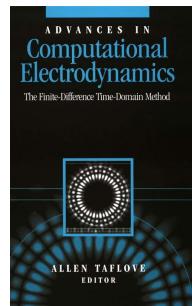
FDTD Solution of the First Two 1-D Scalar Maxwell's Equations / FDTD-Lösung der ersten beiden 1D skalaren Maxwell-Gleichungen



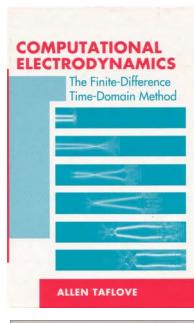
FDTD Books / FDTD-Bücher



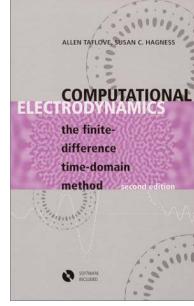
Kunz, K. S., Luebbers, R. J.: *The Finite Difference Time Domain Method for Electromagnetics*. 1993



Taflove, A. (Editor):
Advances in Computational Electrodynamics: The Finite-Difference Time-Domain Method. Artech House, 1998.

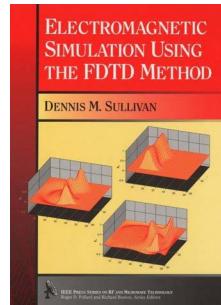


Taflove, A. (Editor):
Computational Electrodynamics: The Finite-Difference Time-Domain Method. Artech House, Boston, 1995.



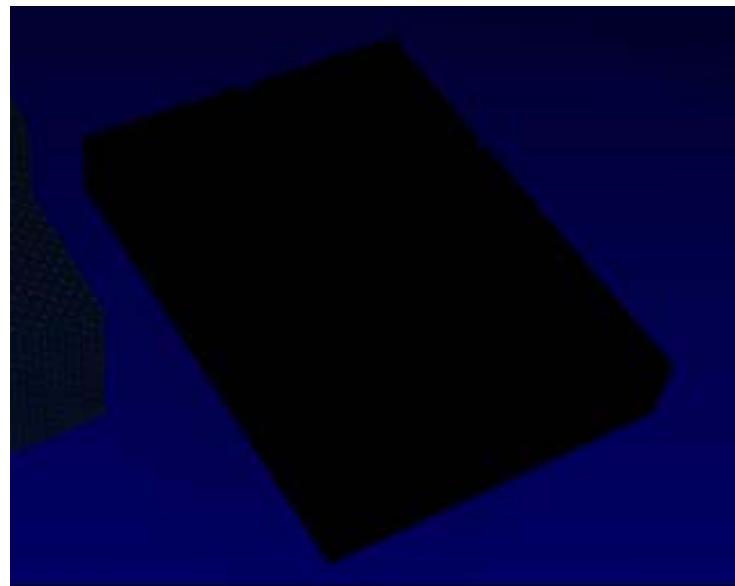
Taflove, A. (Editor):
Computational Electrodynamics: The Finite-Difference Time-Domain Method. 2nd Edition, Artech House, Boston, 2000.

FDTD Books / FDTD-Bücher

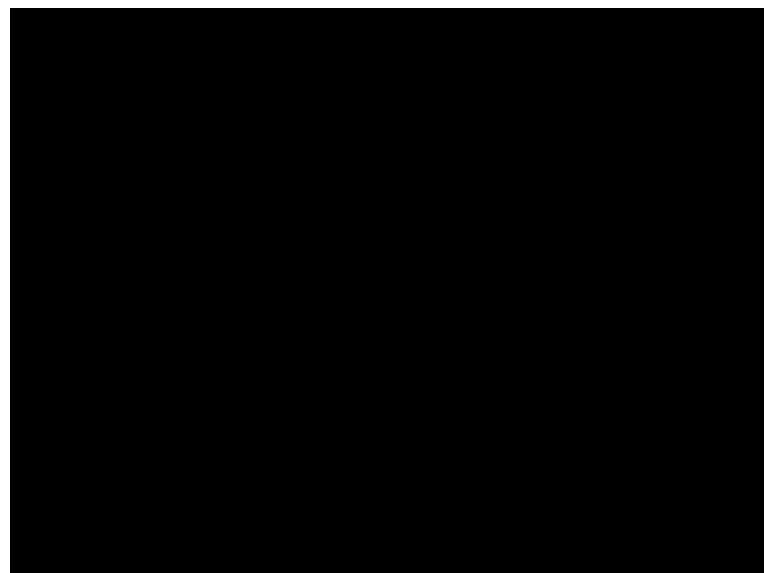


Sullivan, D. M.:
Electromagnetic Simulation Using the FDTD Method. IEEE Press, New York, 2000.

2-D TM FDTD – Photonic Crystals /
2D-TM-FDTD – Photonische Kristalle



2-D TM FDTD – Photonic Crystals /
2D-TM-FDTD – Photonische Kristalle



**End of Lecture 5 /
Ende der 5. Vorlesung**